

Extending computational complexity theory to include thermodynamic resource costs

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with

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SANTA FE INSTITUTE



A continuous-time Markov chain sending $p(0)$ to $p(1) = \sum_j P(i | j) p_j(0)$:

$$\frac{dp_i(t)}{dt} = \sum_j K_{ij}(t)p_j(t)$$

- Example: Dynamics of a digital gate in a circuit
- Example: Dynamics of an entire digital circuit
- Example: Dynamics of a deterministic finite automaton (DFA)
- Example: Dynamics of a Turing Machine (TM)

A continuous-time Markov chain sending $p(0)$ to $p(1) = \sum_j P(i | j) p_j(0)$:

$$\frac{dp_i(t)}{dt} = \sum_j K_{ij}(t) p_j(t)$$

$$\frac{dS(p(t))}{dt} = \dot{Q}(t) + \dot{\Sigma}(t)$$

$$\begin{aligned} \cdot \quad \dot{Q}(t) &= \sum_{ij} K_{ij}(t) p_j(t) \ln \frac{K_{ji}(t)}{K_{ij}(t)} && \text{Entropy flow rate} \\ \cdot \quad \dot{\Sigma}(t) &= \sum_{ij} K_{ij}(t) p_j(t) \ln \frac{K_{ij}(t) p_j(t)}{K_{ji}(t) p_i(t)} && \text{Entropy production rate} \end{aligned}$$

- Entropy production (EP) rate is non-negative

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$$\frac{dS(p(t))}{dt} = \dot{Q}(t) + \dot{\Sigma}(t)$$

Integrate over time: $-\Delta Q = \Delta \Sigma - \Delta S$

- $\Delta S = S(p_1) - S(p_0)$ is gain in **Shannon entropy** of p
- $-\Delta Q$ is (Shannon) **entropy flow** from system between $t = 0$ and $t = 1$
- $\Delta \Sigma$ is total **entropy production** in system between $t = 0$ and $t = 1$
 - **cannot be negative**
(I.e., the second law of thermodynamics)

For many *non-Markovian* chains sending $p(0)$ to $p(1) = \sum_j P(i | j) p_j(0)$:

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GENERALIZED LANDAUER BOUND

- System connected to multiple reservoirs, e.g., heat baths at different temperatures. (So “ kBT ” not defined.)
- **Arbitrary** number of states
- **Arbitrary** initial distribution p_0
- **Arbitrary** dynamics $P(x_1 | x_0)$

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$$-\Delta Q = \Delta \Sigma + S(p_0) - S(p_1)$$

Entropy Production ($\Delta \Sigma$) is non-negative. So:

“*Generalized Landauer’s bound*”

$$\underline{-\Delta Q \geq S(p_0) - S(p_1)}$$

BEYOND GENERALIZED LANDAUER

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- Yes.

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Entropy Production ($\Delta\Sigma$) is non-negative.

*Are there broadly applicable non-negative lower bounds on $\Delta\Sigma$,
to add to the lower bound $-\Delta Q \geq S(p_0) - S(p_1)$?*

- Yes.

- Focus on two: ***Speed limit theorem*** (SLT) and ***Mismatch cost***

Use them to investigate
the (thermodynamic) resource costs of
computational machines

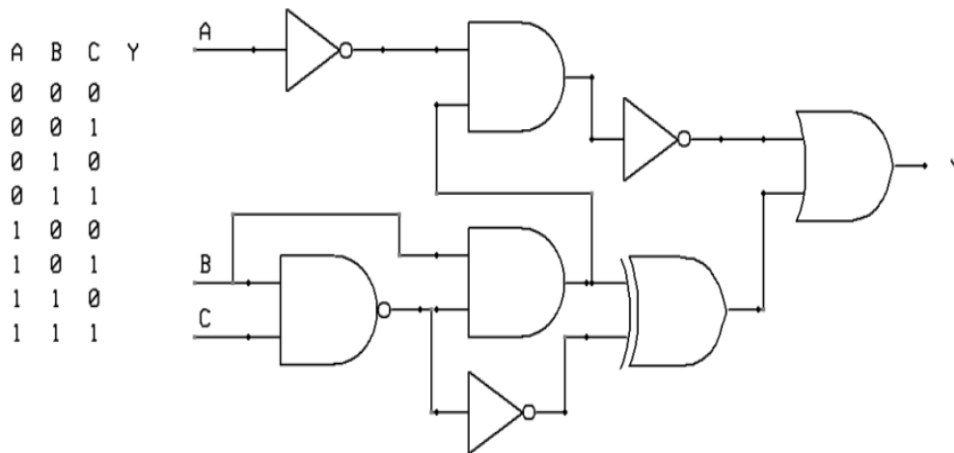
BOOLEAN CIRCUITS

- Currently, all mass-produced computers are implemented with circuits.
- The simplest circuit is one without loops or branches (a “straight-line program”)
- If set of allowed gates are a universal basis (e.g., NAND gates), then can build a circuit with them to implement any desired Boolean function.



$$-\Delta Q = \Delta \Sigma + S(p_0) - S(p_1)$$

- For fixed $P(x_1 \mid x_0)$, changing p_0 changes Landauer cost, $S(p_0) - S(p_1)$
- N.b., the same $P(x_1 \mid x_0)$ - e.g., same AND gate - has different p_0 , depending on where it is in a circuit.
- So even for a thermo. reversible gate ($\Delta \Sigma(p_0) = 0$), **changing the gate's location in a circuit** (changes $S(p_0) - S(p_1)$ and so) **changes $-\Delta Q(p_0)$**

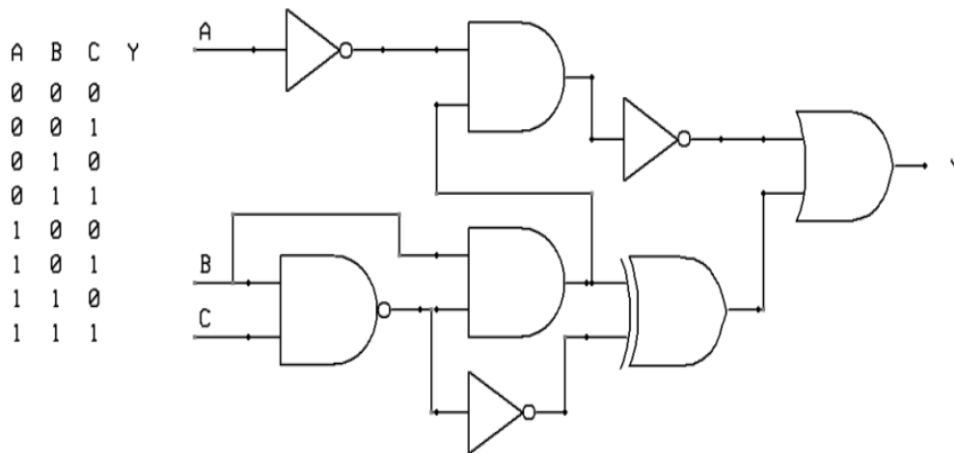


- Changing a gate's location in a circuit changes $S(p_0) - S(p_1)$, and so changes the heat it produces, $-\Delta Q(p_0)$
- Sum those heats over all gates to get minimal heat flow of that circuit

Different circuits implementing same Boolean function on same input distribution have different minimal heat

- Formally, those differences in minimal heat of the circuits are differences in EPs of the circuits, arising due to modularity of gates

➤ A new circuit design optimization problem



Demaine, E., et al., *Comm. ACM*, 2016

- Considers a similar problem - but incorrectly sets Landauer cost at each gate to same value, $KT \ln(2)$.

Original speed limit theorem (SLT): $\Delta\Sigma \geq \frac{L(p(0), p(1))}{A_{0,1}}$

- $L(p(0), p(1))$: L_1 distance from distribution $p(0)$ to distribution $p(1)$
- $A_{0,1}$: total number of (stochastic) state jumps from $t = 0$ to $t = 1$

Since introduced, SLT has been strengthened several ways (more complicated formulas).

Shiraishi, N., Funo, K.; Saito, K., *PRL* (2018)

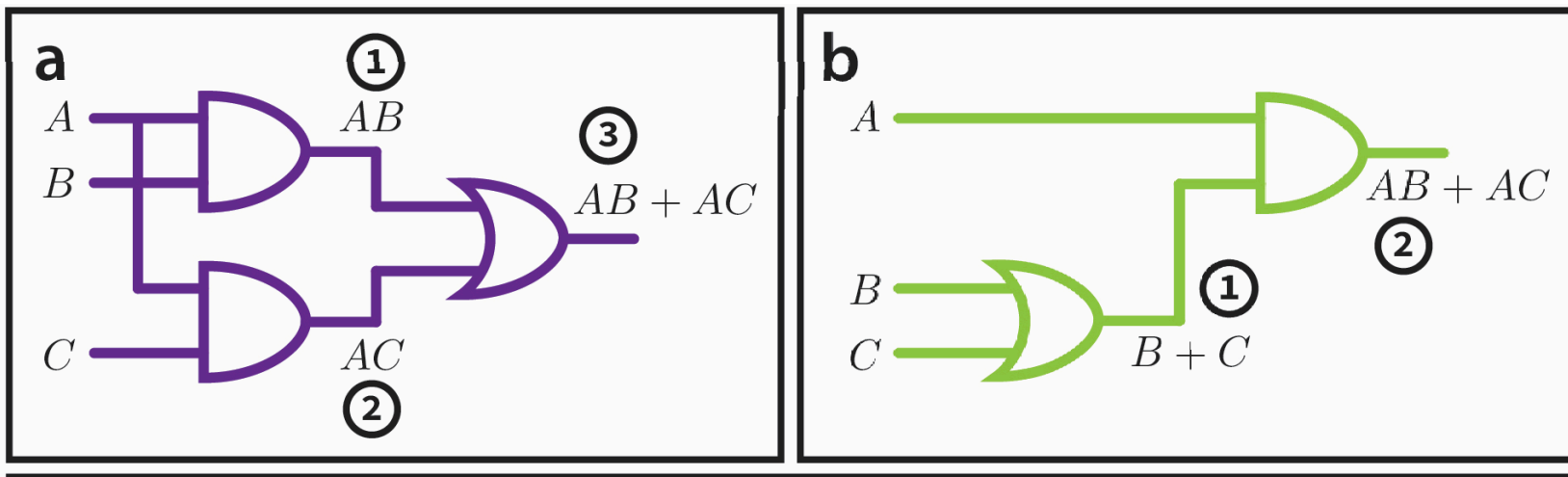
Delvenne, J., Falasco, G.; arXiv:2110.13050

Lee, J., et al.; *PRL* (2022)

Van Vu, T., Saito, K.; *PRL* (2023)

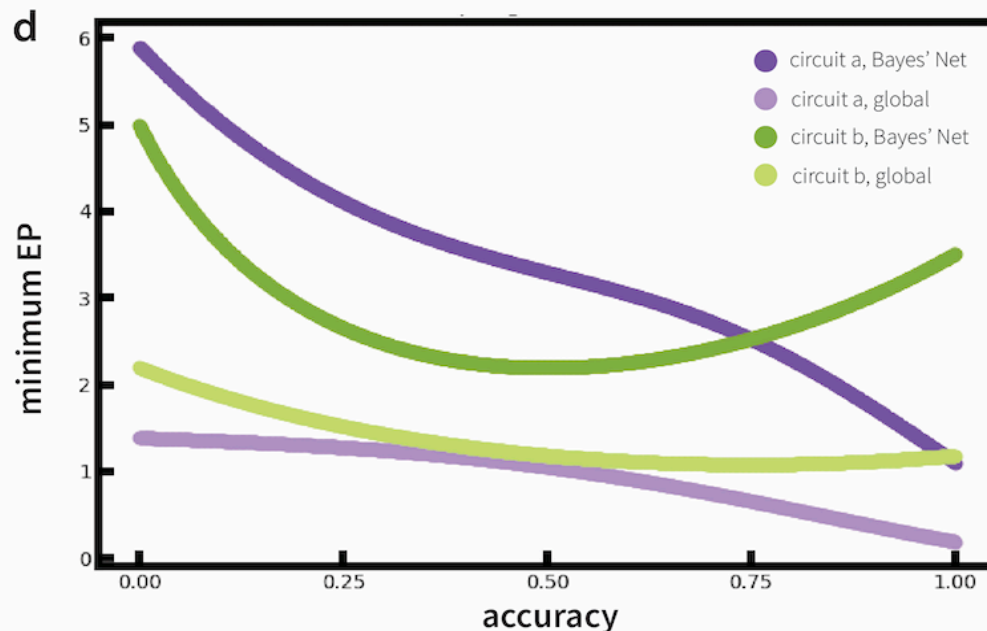
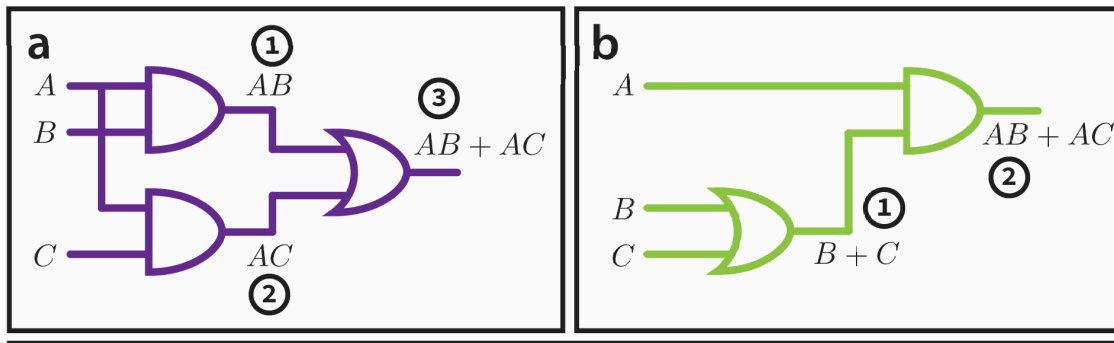
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- $A_{0,1}$: total number of (stochastic) state jumps from $t = 0$ to $t = 1$
- Suppose uniform initial distribution over all gates and input bits;
- How does the (Lee et al.) SLT bound vary with error rate of gates, for two logically equivalent circuits?



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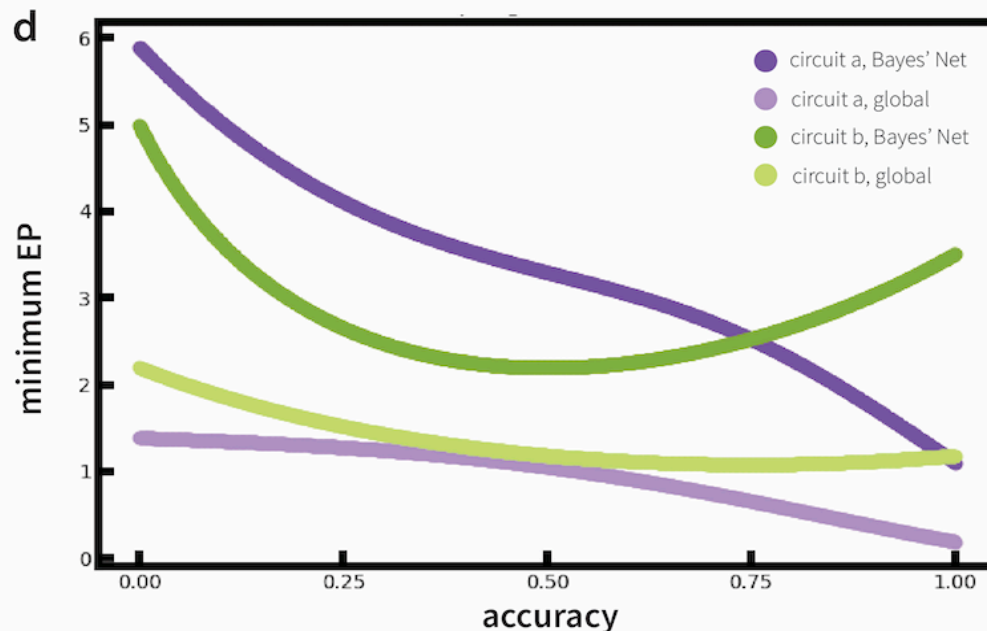
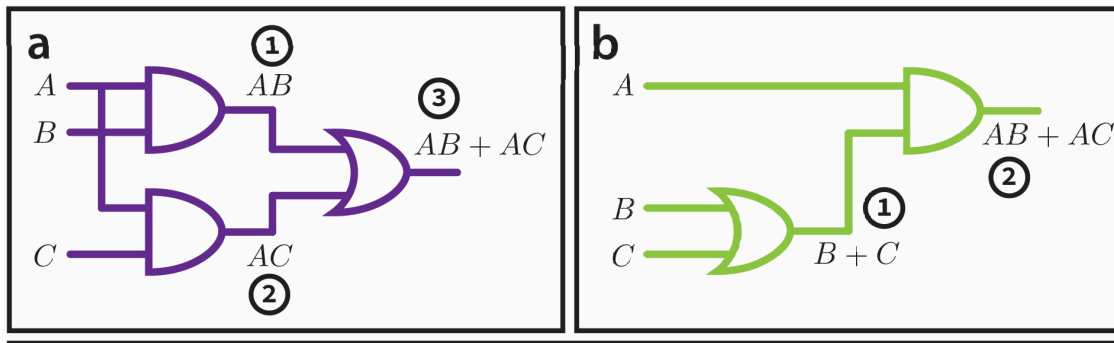
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Tasnim, F., Wolpert, D.,
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(2023)

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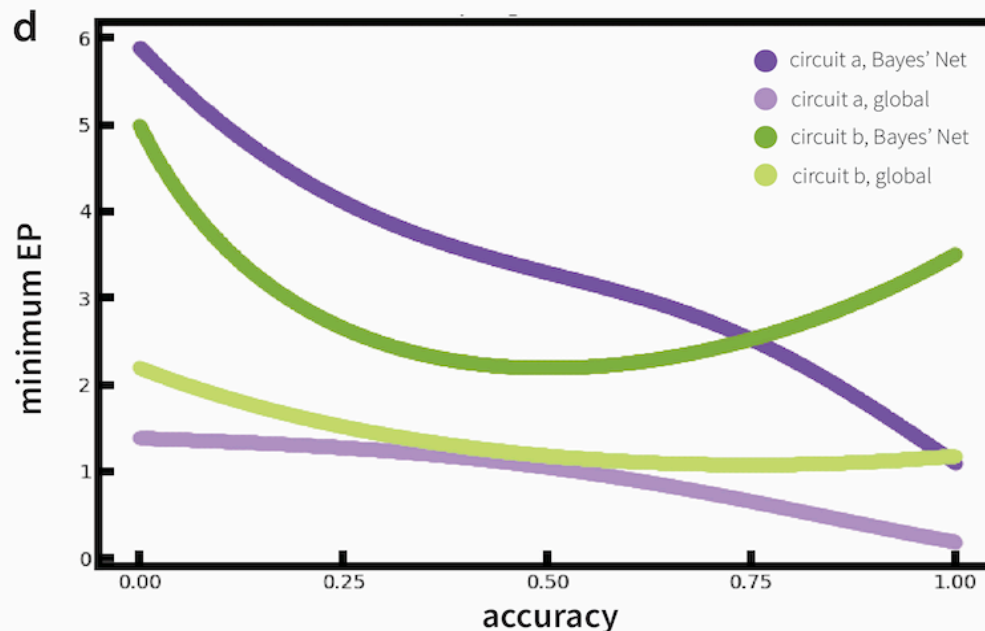
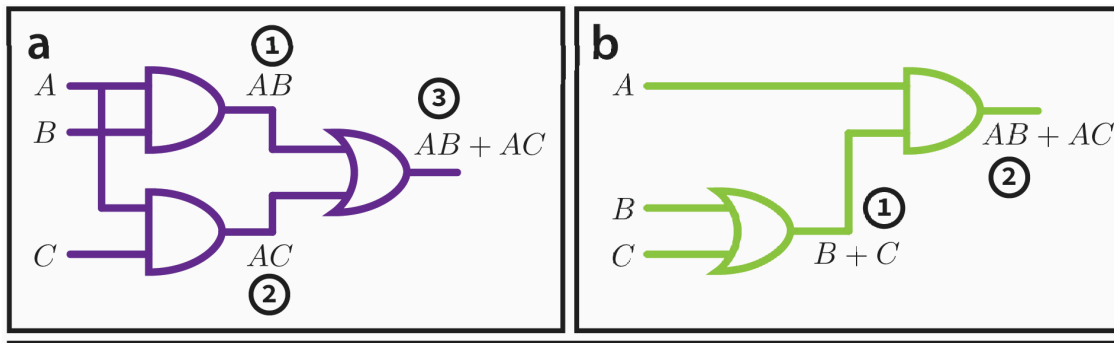
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What are curves for other circuits?

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A: Who knows!

DEPENDENCE OF EP ON INITIAL DISTRIBUTION

- Arbitrary dynamics $P(x_1 | x_0)$
- ***Assume system is thermo. reversible for initial distribution q_0***

I.e., $\Delta\Sigma(q_0) = 0$

- Run that system with initial distribution $p_0 \neq q_0$ instead:

$$\Delta\Sigma(p_0) = D(p_0 || q_0) - D(p_1 || q_1) \geq 0$$

where $D(. || .)$ is relative entropy (KL divergence)

Wolpert, D., Kolchinsky, A., *New J. Phys.* (2020)

Riechers, P., Gu, M., *Phys. Rev. E* (2021)

Kolchinsky, A., Wolpert D., *arxiv:2103.05734*

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Any nontrivial process that is thermodynamically reversible for one initial distribution ***will be costly*** for any other initial distribution

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where $D(. \parallel .)$ is relative entropy (KL divergence)

$D(p_0 \parallel q_0) - D(p_1 \parallel q_1)$ is called ***mismatch cost***

Mismatch cost example/

- **Two** distinct bit-erasing gates, each with thermo. rev. initial distribution q_0
- Run gates in parallel, on bits x^A and x^B , with initial distribution $p_0(x^A, x^B)$
- Assume $p_0(x^A) = q_0(x^A)$ and $p_0(x^B) = q_0(x^B)$.
- So each gate, by itself, generates zero EP. But:

*If $p_0(x^A, x^B)$ statistically couples the bits, then
full system is **not** thermo. reversible,
and **generates nonzero EP***

- **Formally**: Since gates are distinct, the thermo. rev. *joint* distribution is $q_0(x^A, x^B) = q_0(x^A)q_0(x^B)$.

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- **Intuition:** Running two thermo. reversible gates in parallel loses information in their initial coupling, and so is not thermo. reversible.

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- **Broader lesson:** *Modularity increases EP*

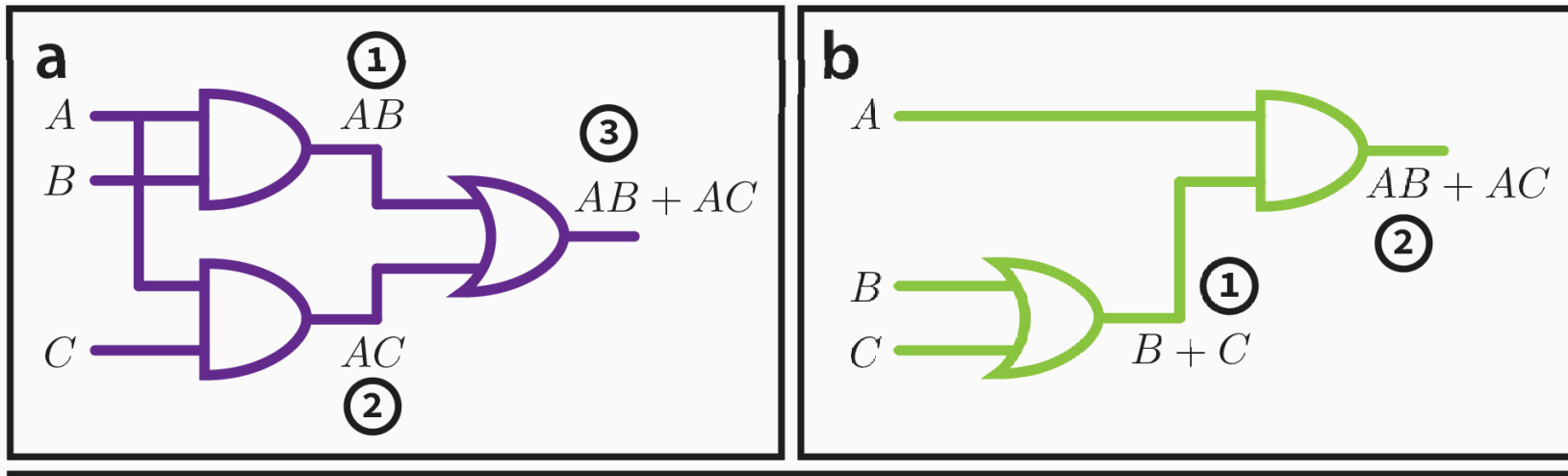
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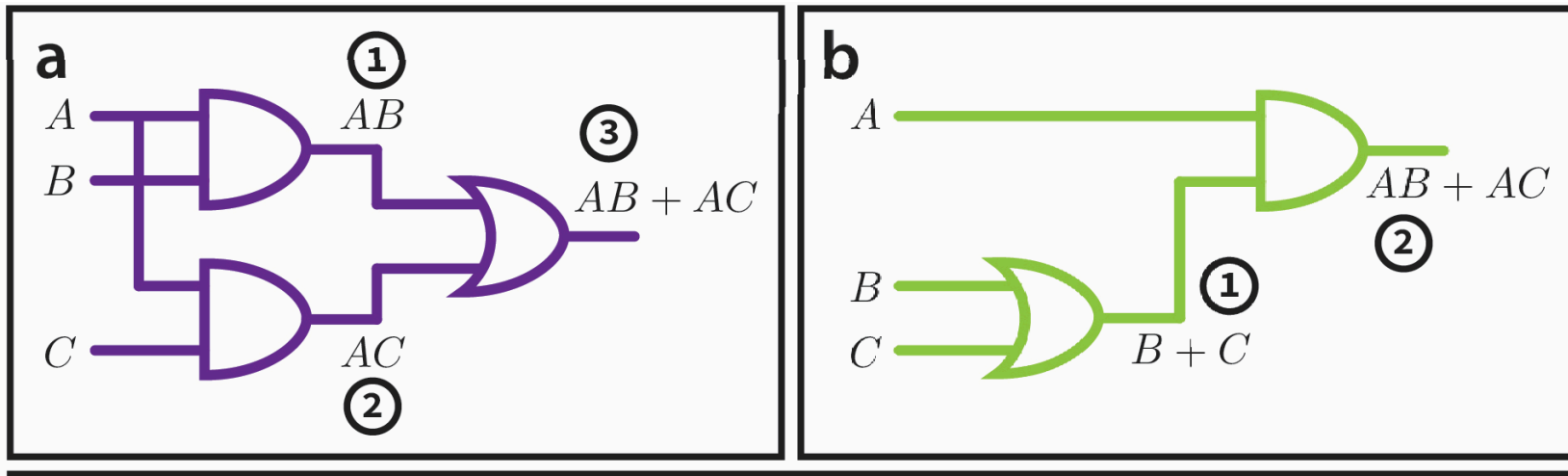
- **Broader lesson:** Whatever its practical benefits might be,
modularity is thermodynamically costly (!)

MISMATCH COST OF BOOLEAN CIRCUITS



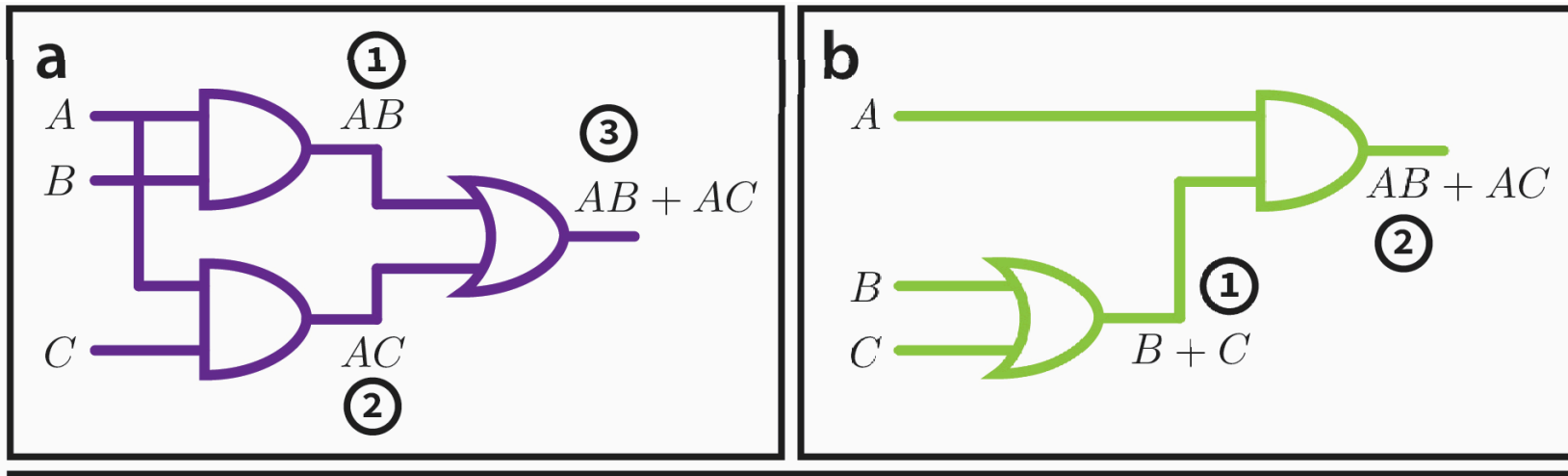
- Physical process updating each gate in a real circuit depends only on that gate's inputs – it is independent of all other gates.
- Similar to parallel bit erasure.

MISMATCH COST OF BOOLEAN CIRCUITS



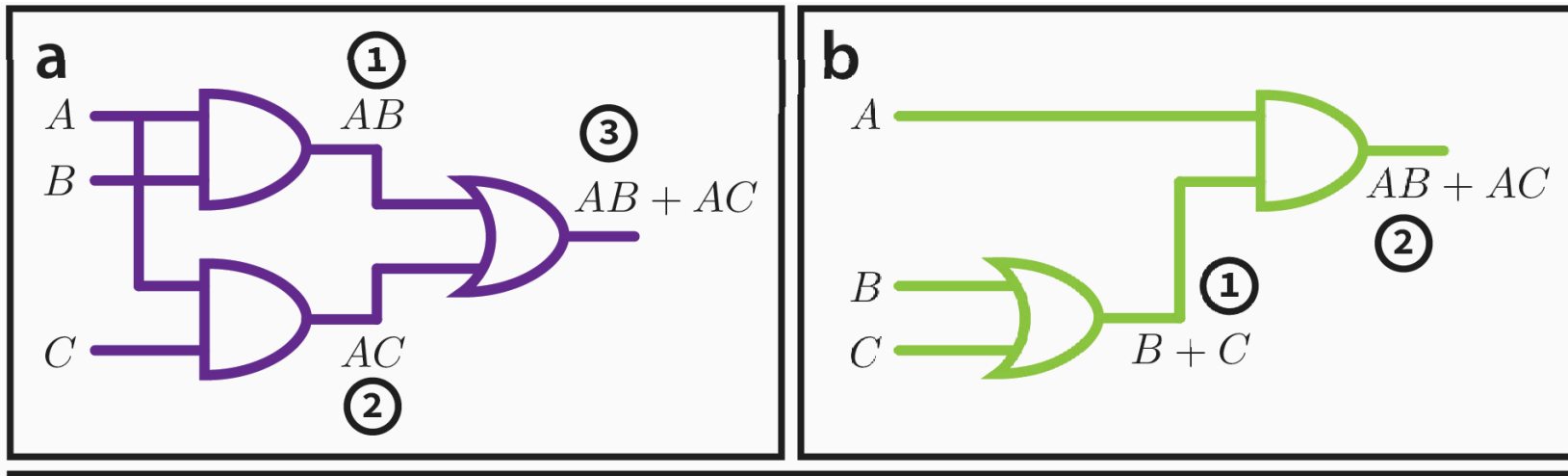
- On first use of circuit, inputs and all gates uniformly random
 - Assume priors of gates are also all uniformly random
-
- Then mismatch cost = 0 - for the first use of the circuit

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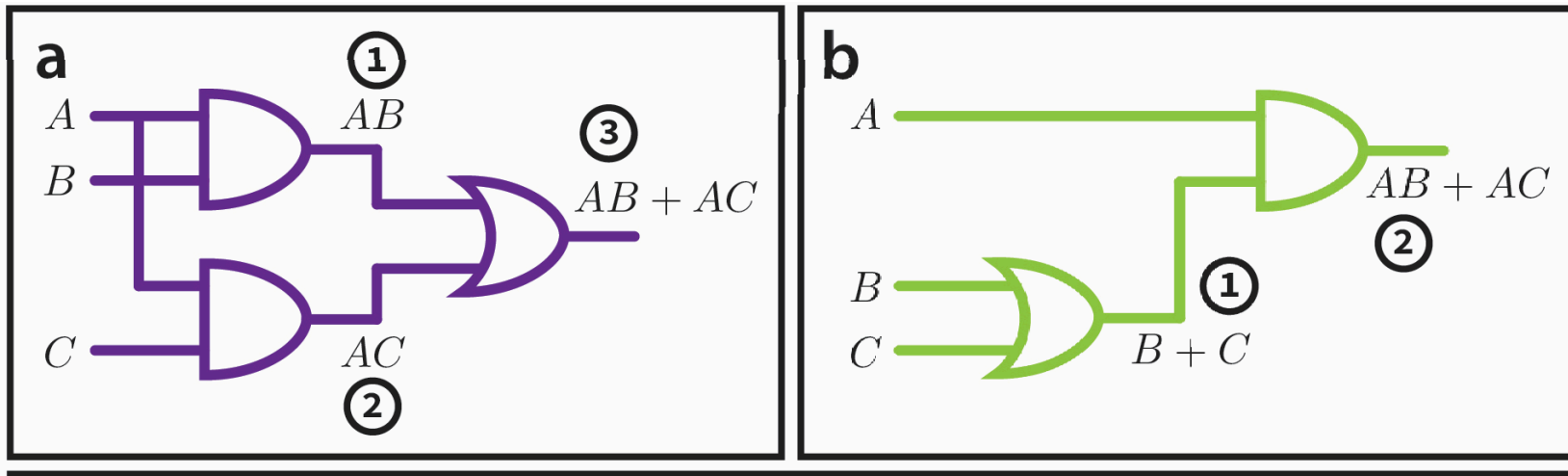
- On first use of circuit, inputs and all gates uniformly random
- Assume priors of gates are also all uniformly random
- Suppose on second use, inputs are again uniformly random – but gates are reinitialized, e.g., to uniformly random.
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MISMATCH COST OF BOOLEAN CIRCUITS



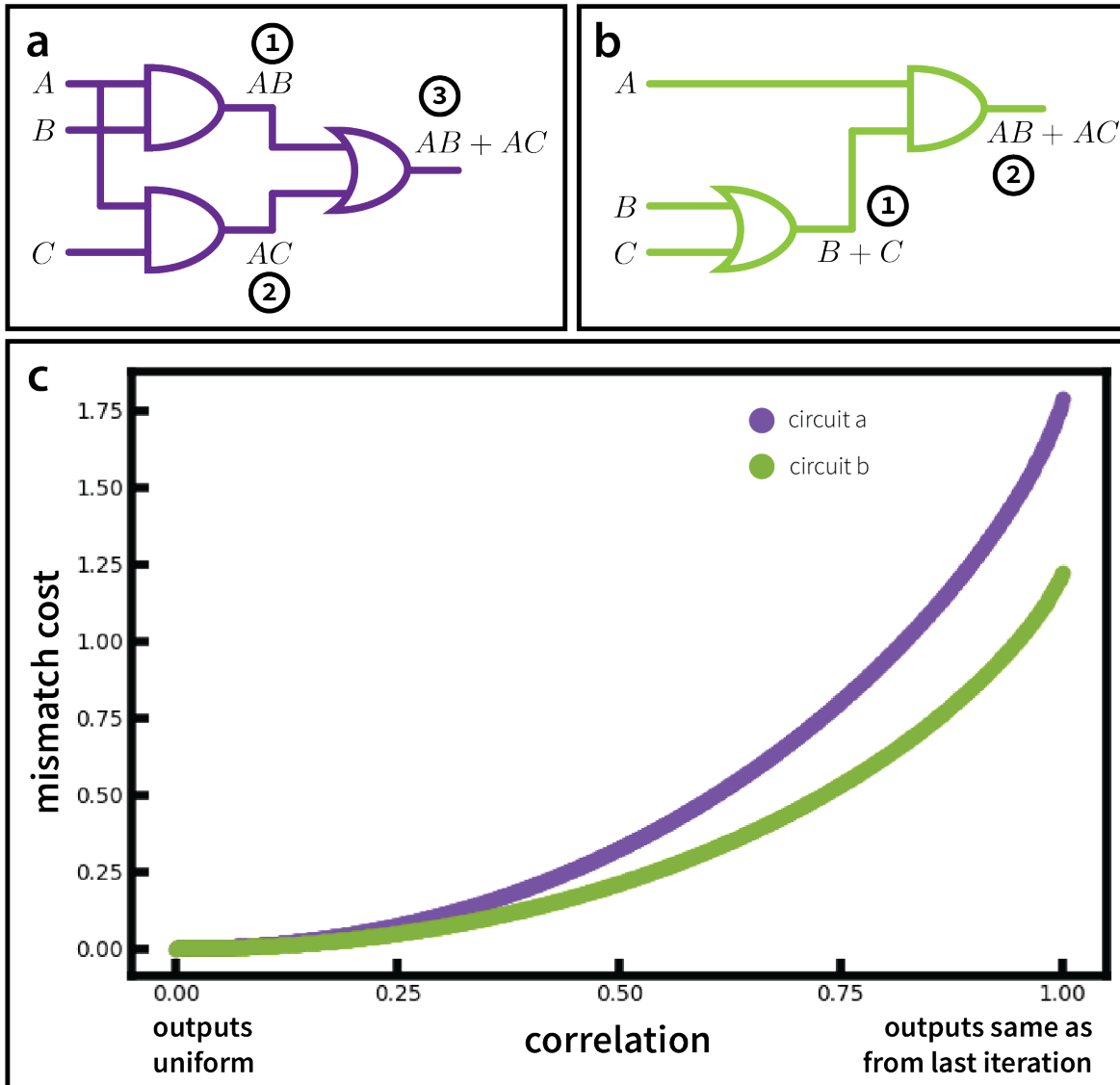
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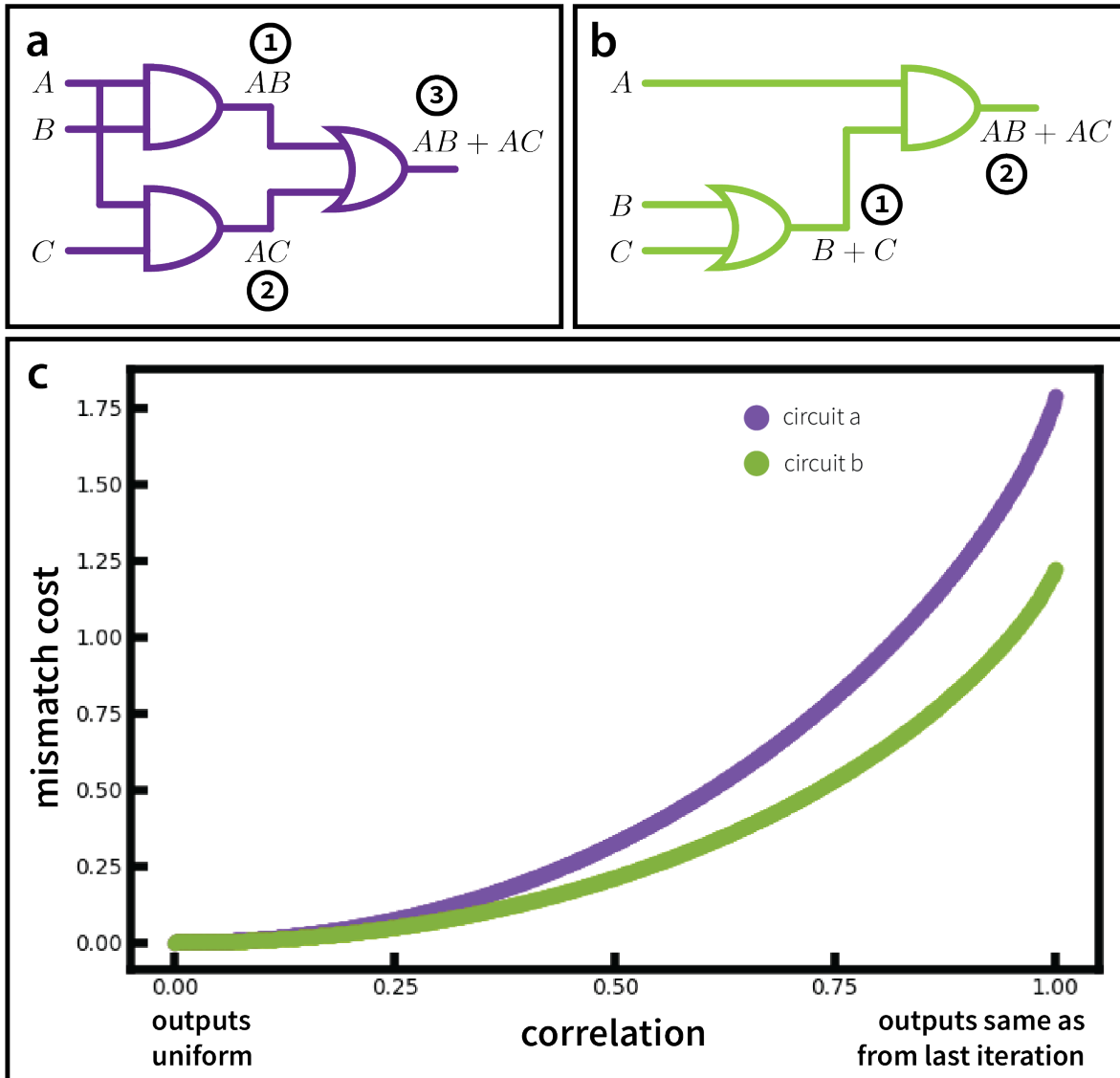
- On first use of circuit, inputs and all gates uniformly random
- Assume priors of gates are also all uniformly random
- Suppose on second use, inputs are again uniformly random – *but gates still have their values from end of first use.*
- **Then mismatch cost $\neq 0$** - for the second use of the circuit.

MISMATCH COST OF BOOLEAN CIRCUITS



Tasnim, F., Wolpert, D.,
Korbel J., Lynn, C., et al.
(2023)

MISMATCH COST OF BOOLEAN CIRCUITS



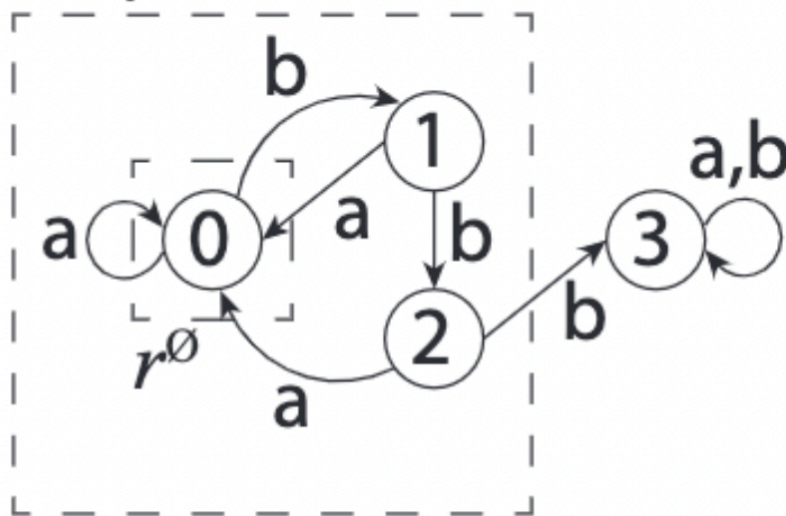
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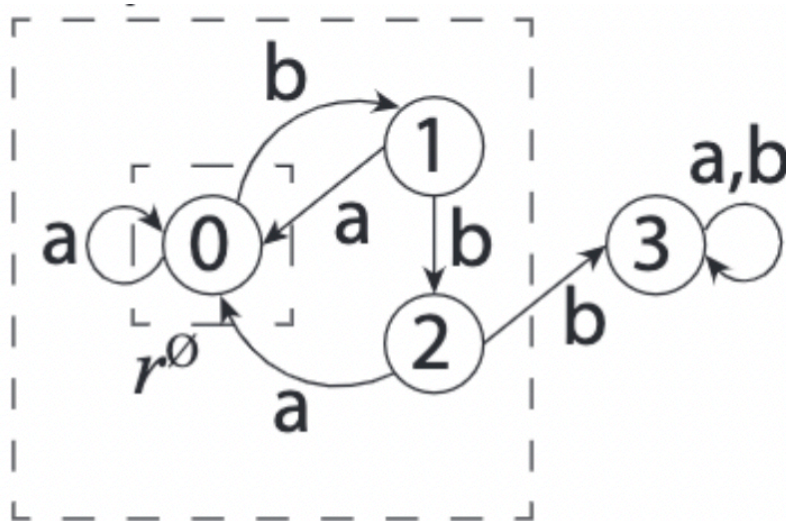
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DETERMINISTIC FINITE AUTOMATA (DFA)

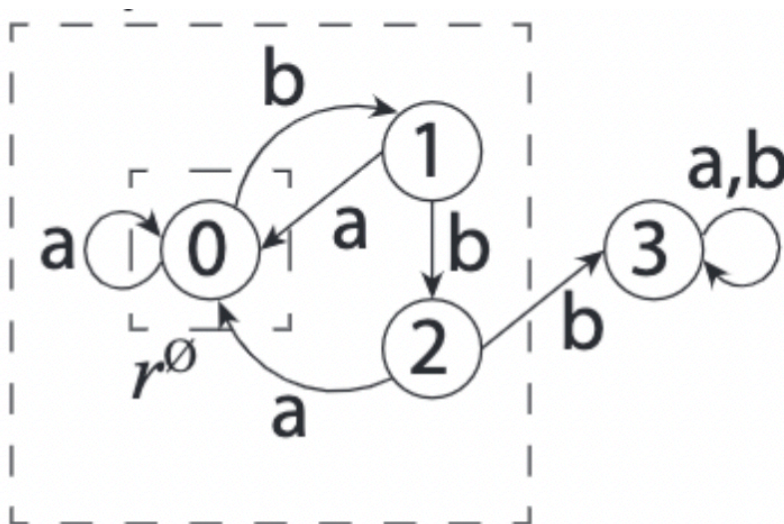
- **Simplest computational machine in Chomsky hierarchy**
 - Finite number of states; one **initial state**, multiple “**accept states**”
 - Feed in a finite string of bits;
 - Each (bit, state) pair maps to a new state, after which next bit is read
 - A DFA “**accepts**” a string if it causes the DFA to end in an accept state
 - “**Language**” of a DFA is all input strings that it accepts
 - Many languages that are not accepted by any DFA
- **Example:** DFA that accepts any string with no more than two successive ‘b’ bits:



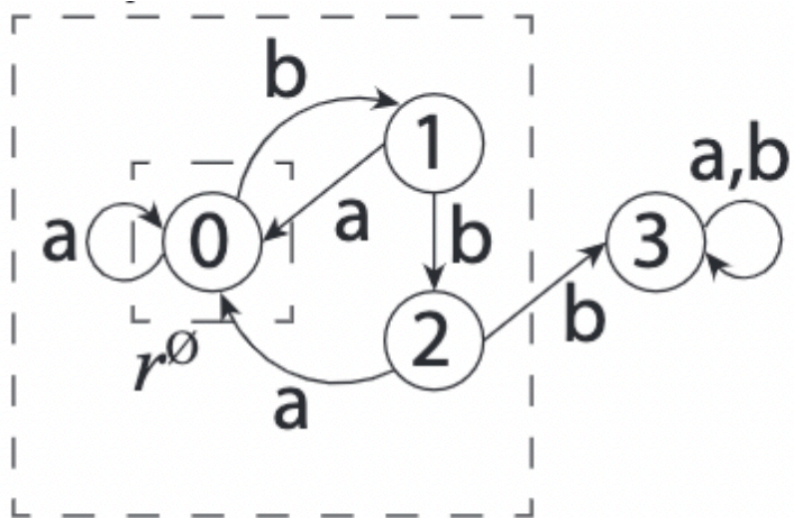
- Every digital computer is “local”
 - the only part of memory any processing unit is directly physically coupled to is its current input
- E.g., in a DFA, state update only physically coupled to current input symbol, not any earlier / later symbols
- Results in “modularity (mismatch) cost” – just like parallel bit erasure



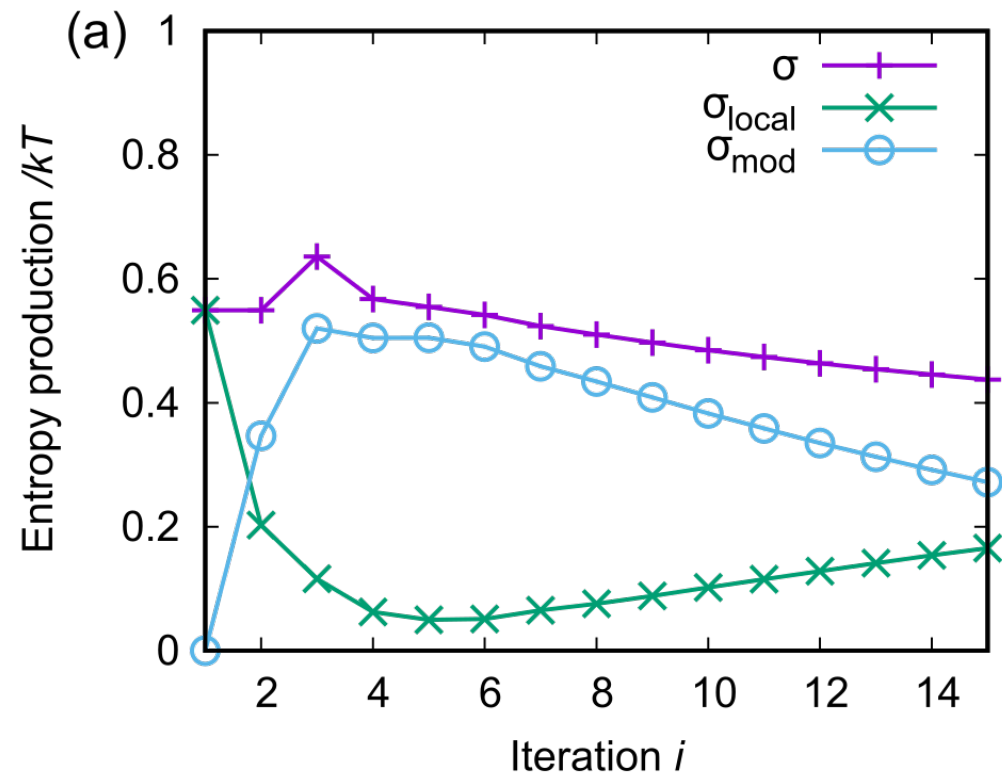
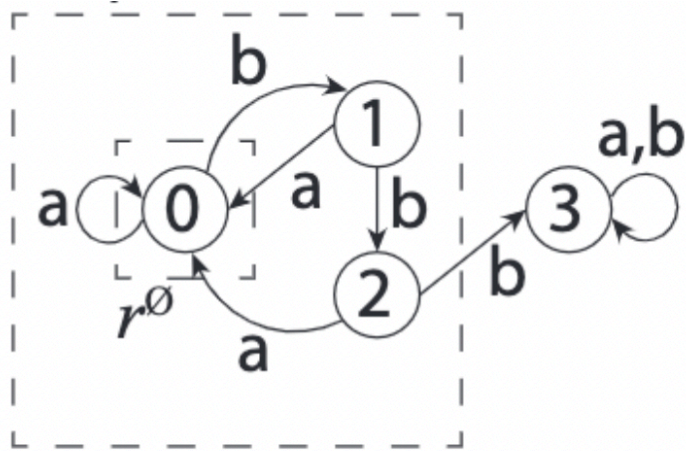
- Every (synchronous) digital computer is “periodic”
 - every successive iteration is the same physical process, and so in particular has the same prior.
- E.g., in a DFA, every iteration has same prior
- So if prior = actual distribution for iteration i (so zero mismatch cost), they will differ for iteration $i + 1$ in general (so nonzero mismatch cost!)
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- Total mismatch cost = modularity cost + local cost

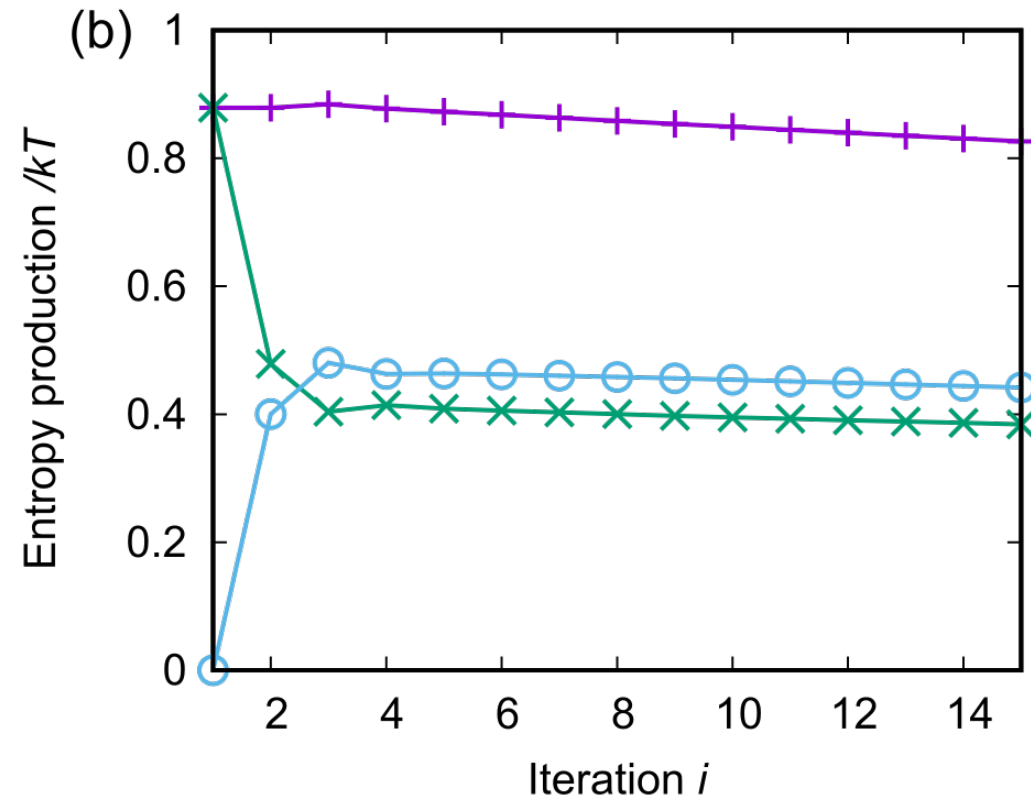
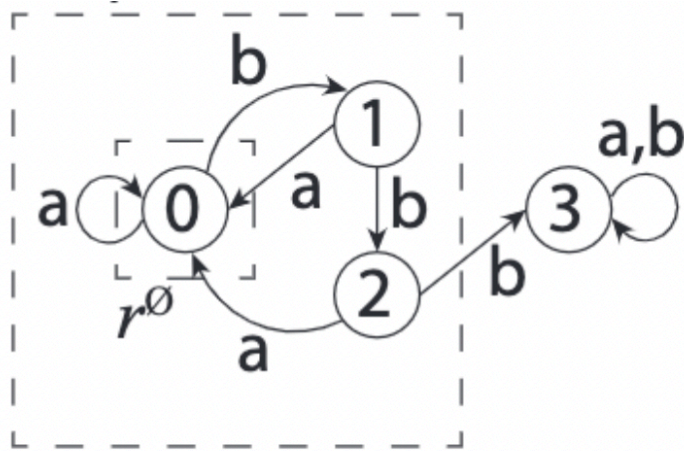


EXAMPLE



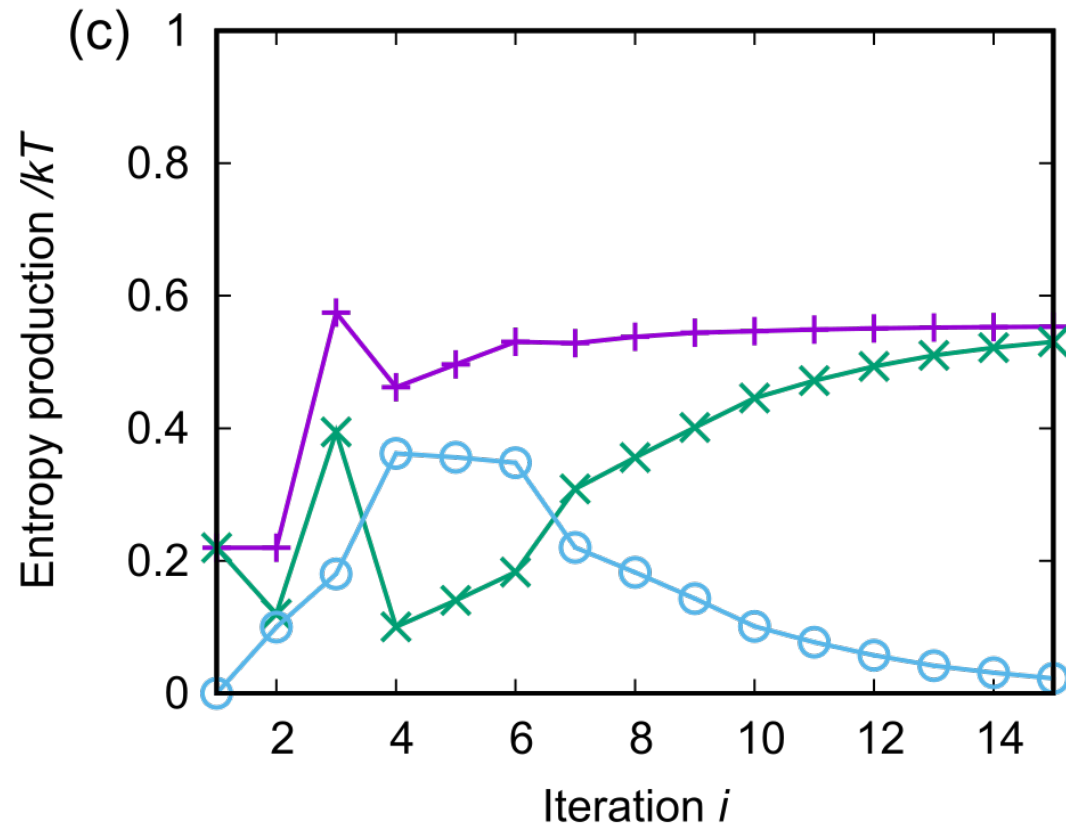
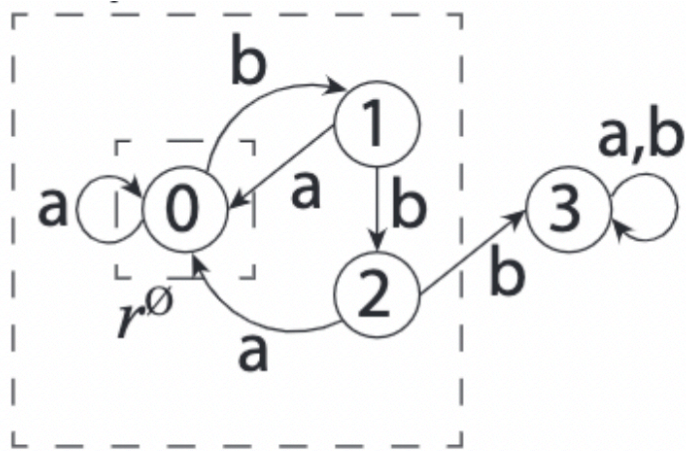
Input strings have IID symbols with equal probability of a and b

EXAMPLE



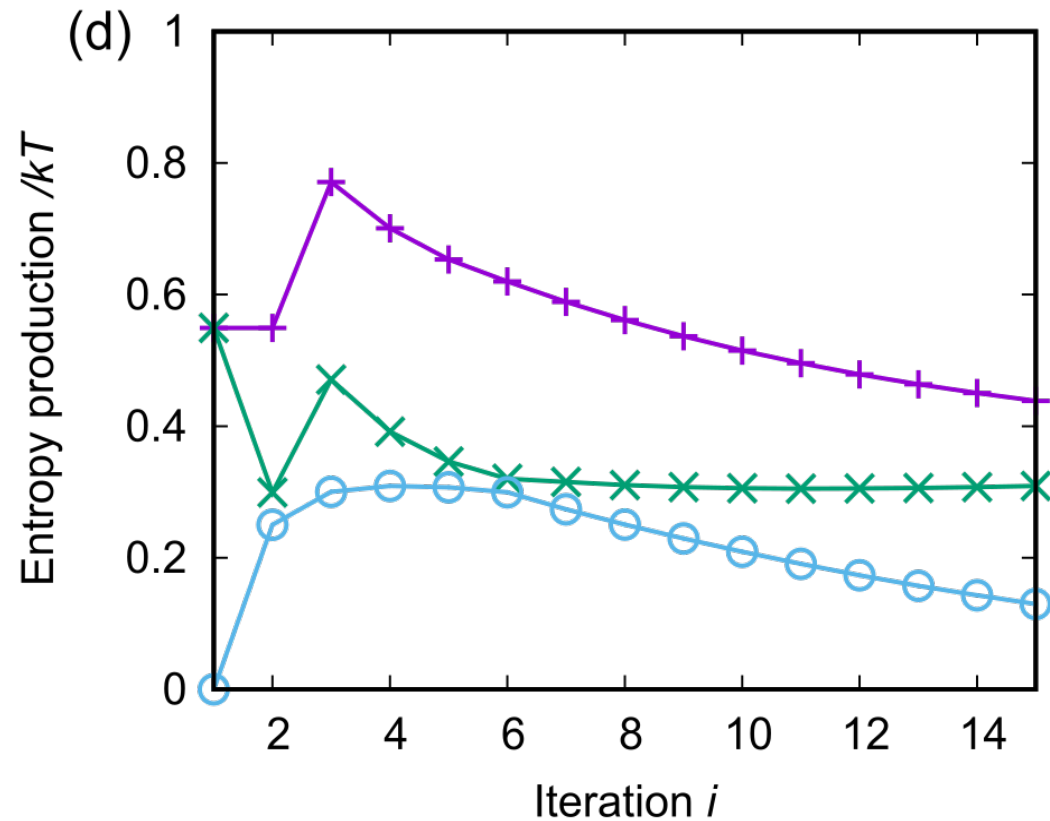
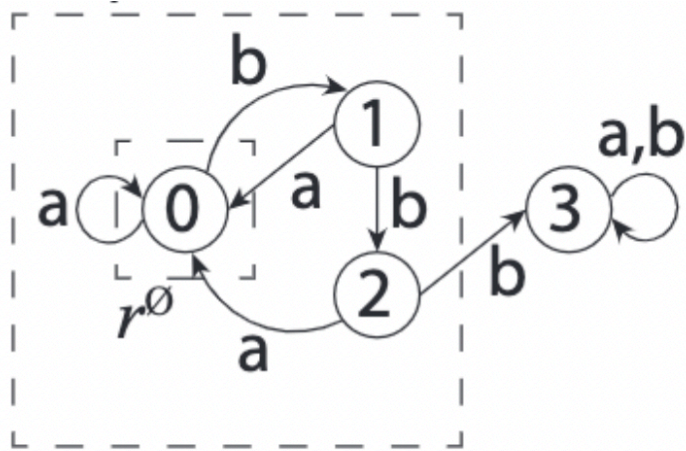
Input strings have IID symbols with probability of $a = 0.8$

EXAMPLE



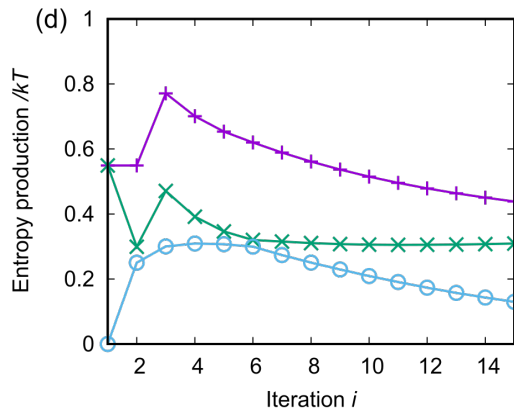
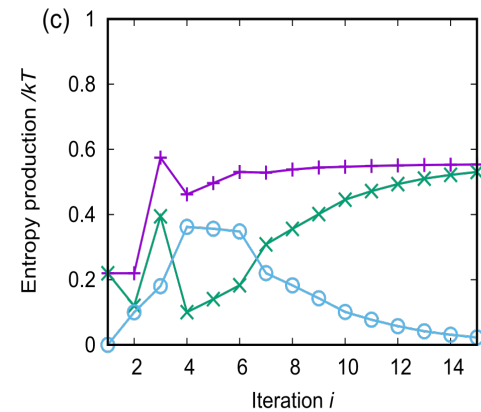
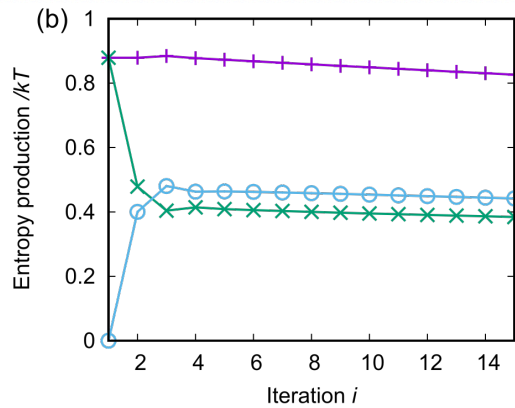
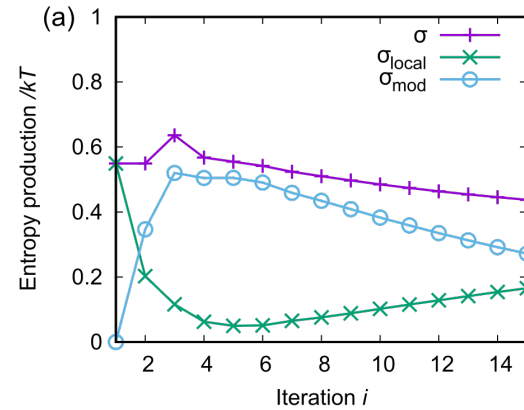
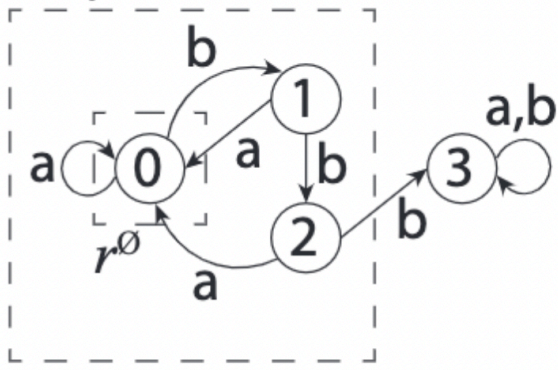
Input strings are first order Markov chains (starting from uniform probability)

EXAMPLE



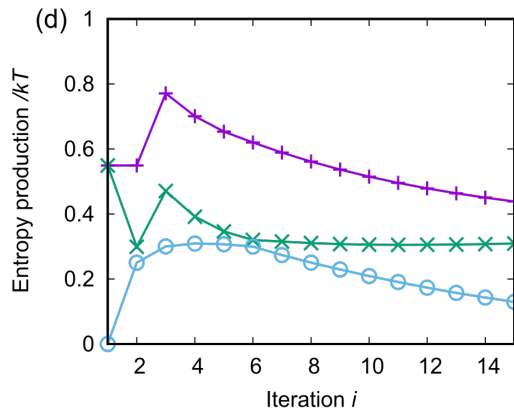
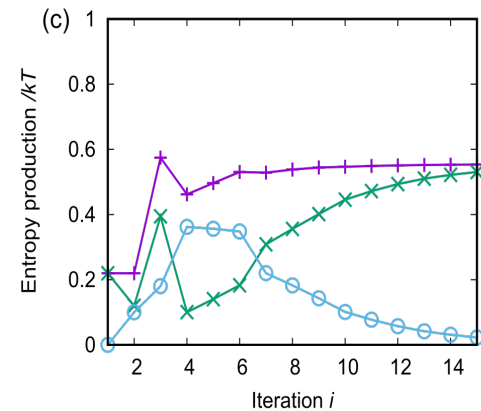
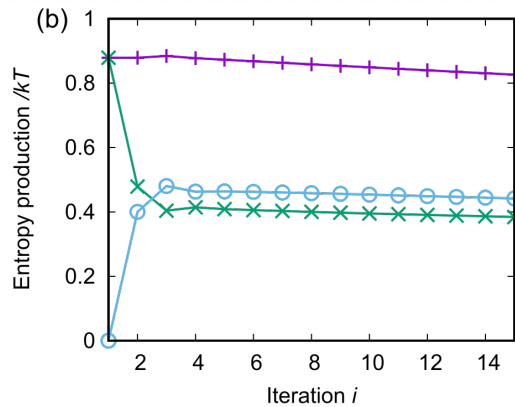
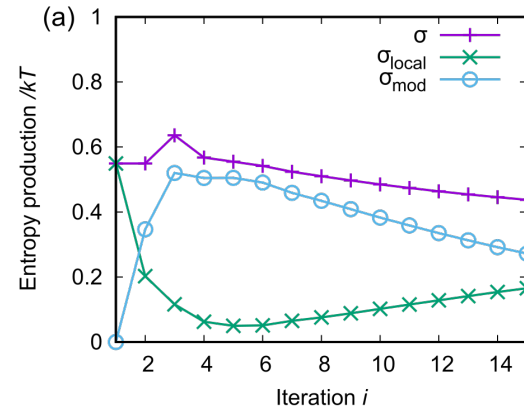
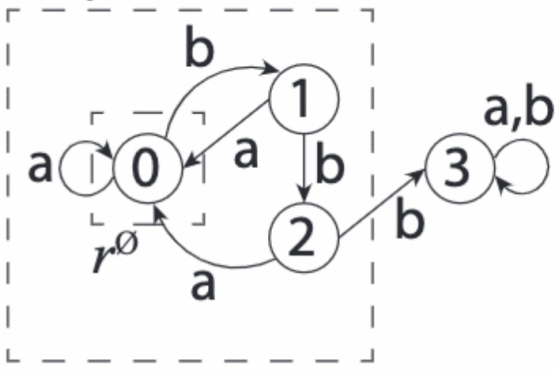
Input strings have IID symbols with probability of $a = 0.8$

EXAMPLE



Ouldrige, T., Wolpert, D., arxiv:2208.06895 (2022)

EXAMPLE



What causes curves to have these shapes?
What are curves for other DFAs?

A: Who knows!

All these bounds on thermodynamic cost of computers hold *independent of nitty gritty details of physical system implementing the computer*

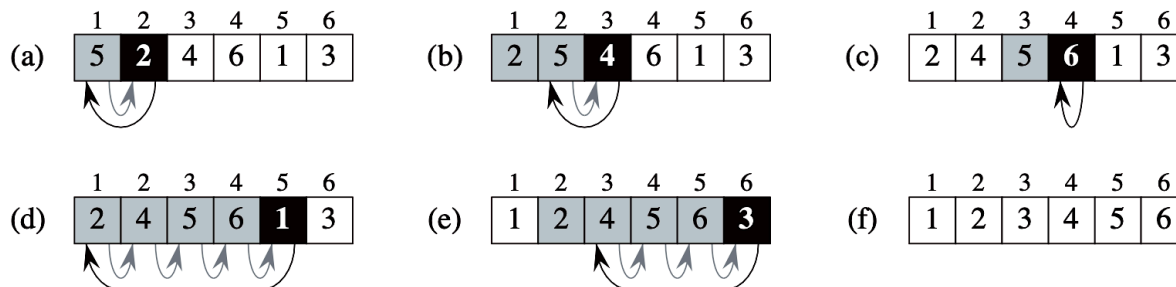
Other recent results based on “nitty gritty details” of computers implemented using CMOS technology.

Less abstract than computational machines – but not deep in the weeds CMOS technology:

Mismatch cost for implementing pseudo-code
(Periodic process, but for simplicity, not local)

- Algorithm to sort any list of six integers into ascending order

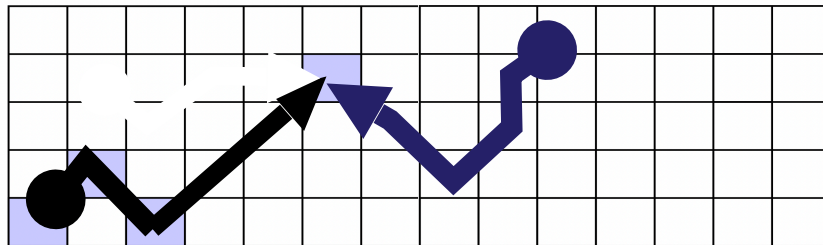
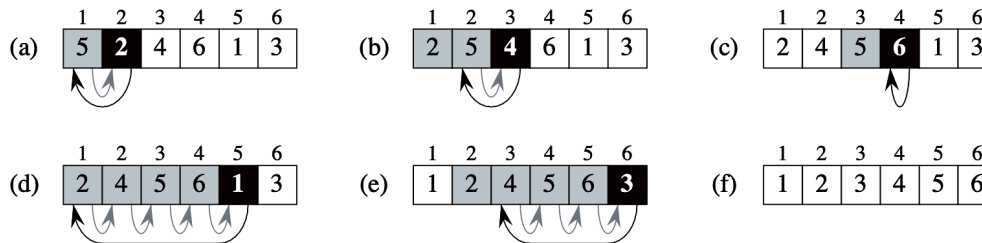
Insertion sort



- All contributions to mismatch cost come from many-to-one maps over the sequence of six integers.

**Just like many-to-one maps cause nonzero “Landauer cost”,
many-to-one maps cause nonzero mismatch cost**

Insertion sort



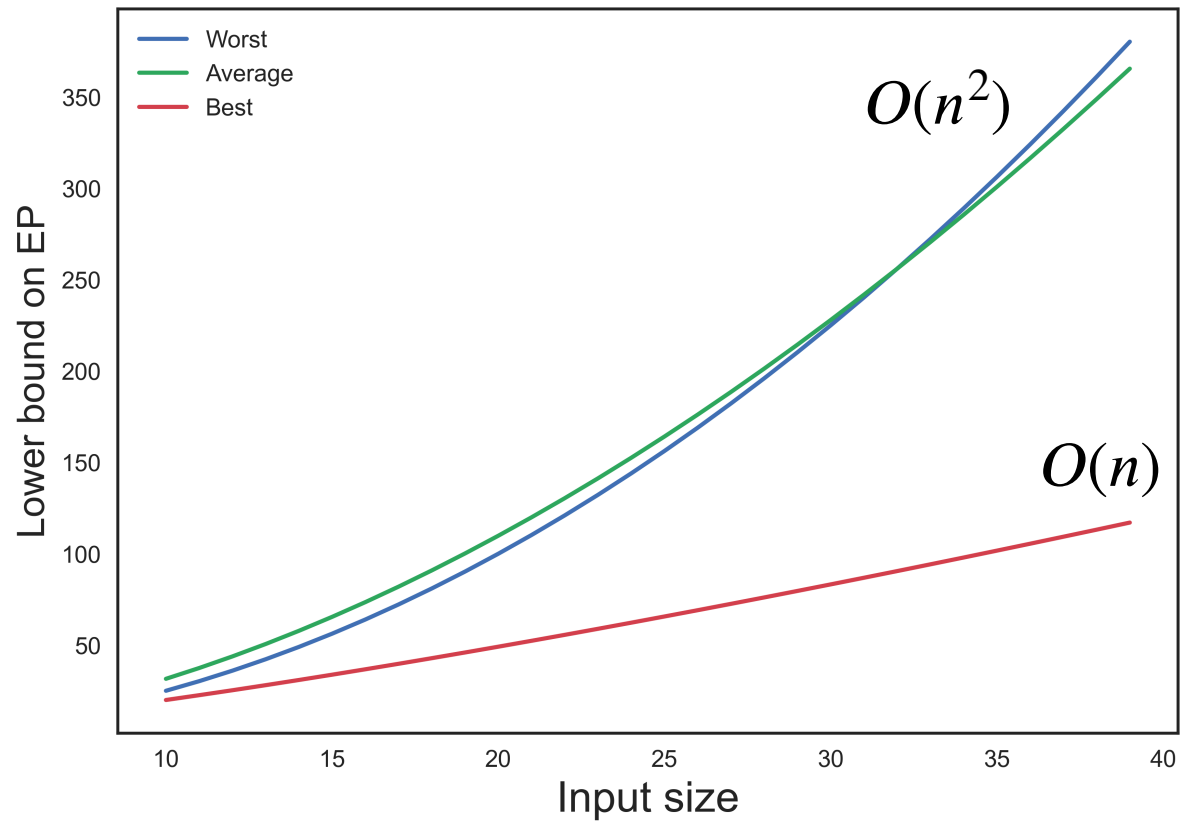
- Each cell: different joint value of variables in the pseudo-code
- Trajectories can merge
- many-to-one maps

Total mismatch cost summed along a trajectory for a “maxent” (uniform) prior:

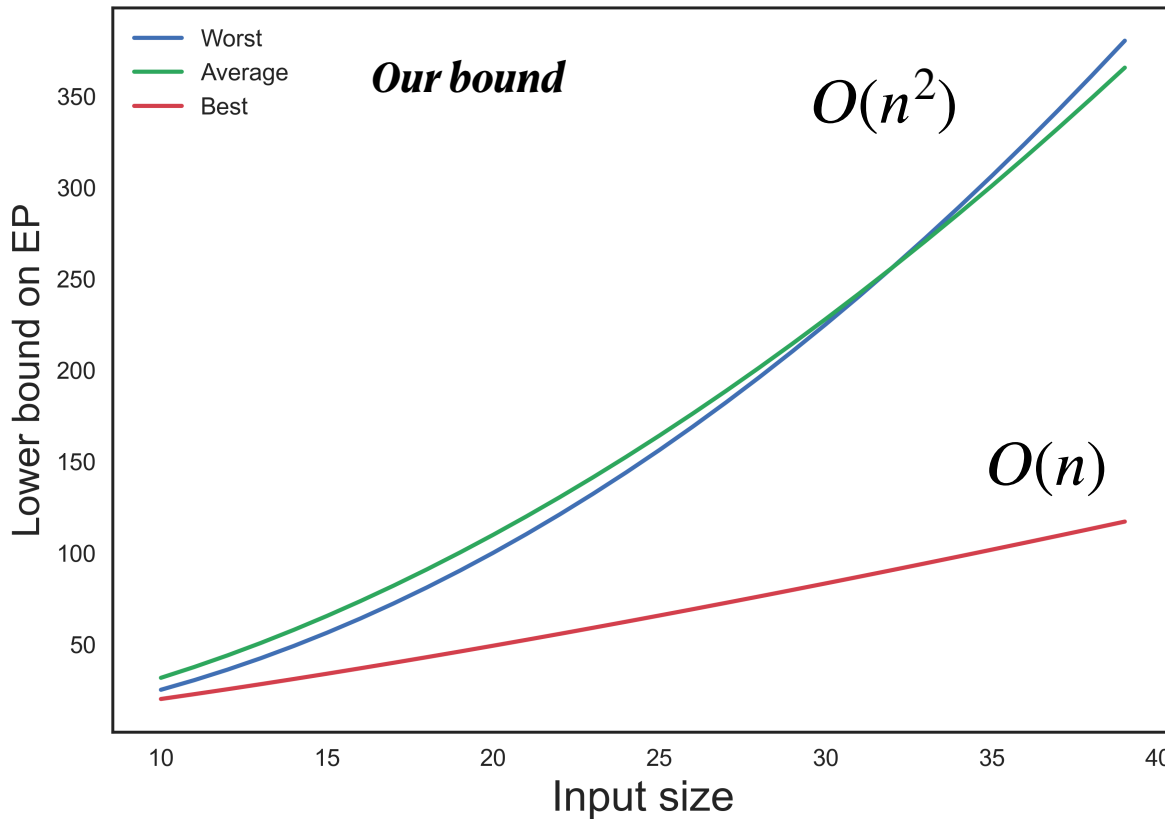
$$\mathcal{A}(\tau) = \ln(\epsilon_{X_1} \cdots \epsilon_{X_\tau}) = \sum_{s=1}^{\tau} \ln \epsilon_{X_s}$$

(where ϵ_{x_i} is “entrance rate” into state x_i , i.e., $\epsilon_{x_i} = \sum_{x_{i-1}} P(x_i | x_{i-1})$)

Energetic complexity!



Energetic complexity!



Landauer's bound

Time	Space (words)	Energy (bits)
$\Theta(n^2)$	$\Theta(n)$	0

Energy-Efficient Algorithms

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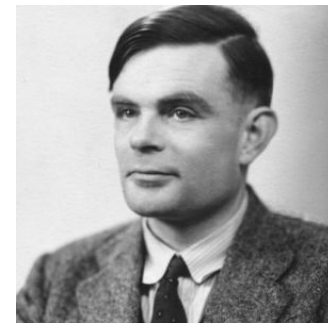
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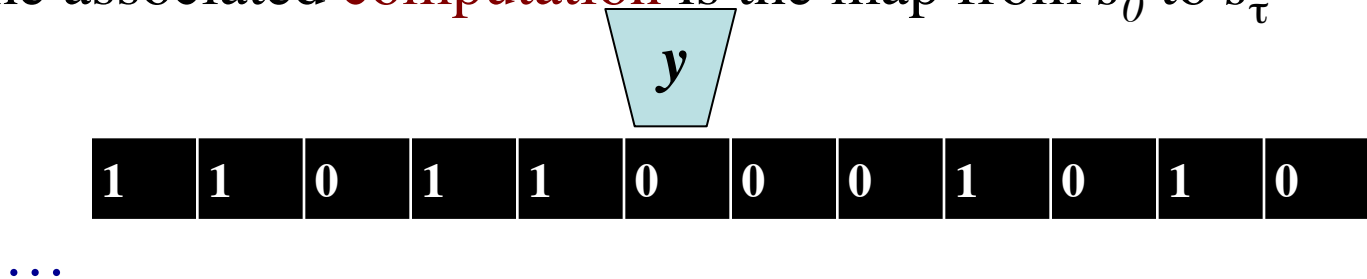
Thermodynamics of Turing Machines

- There are many different abstract models of computers, with different computational powers.
- A particularly important one is the **Turing machine (TM)**
- **Church-Turing thesis:** *“Every function which would naturally be regarded as computable ... is computable by a Turing machine.” (Including computations in the human brain.)*



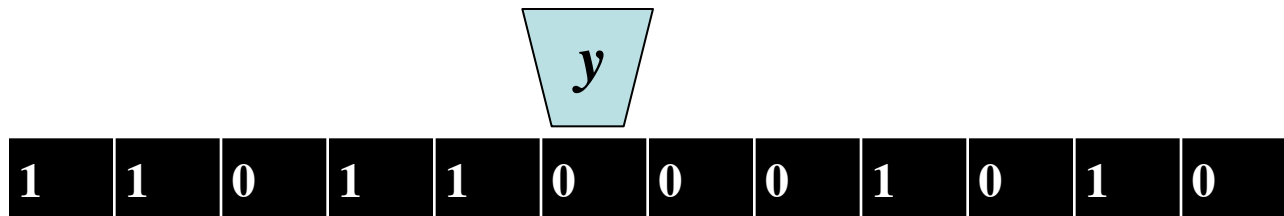
Turing Machines

- 1) Bi-infinite bit string (“tape”) s .
- 2) “**Head**” with n internal states y , one a “**halt state**”
- 3) At each t , head is located at bit b_t which has value $s(b_t)$.
- 4) At each t , based on (y_t, b_t) , the head:
 - i) changes its state to y_{t+1} ;
 - ii) writes a new binary value at b_t ;
 - iii) moves, by up to one bit, in either direction on the tape
- 5) Computation ends (if ever) at time τ if head halts then.
- 6) The associated **computation** is the map from s_0 to s_τ



Turing Machines

- 1) The standard model of computation (up to and including human “computers”)
- 2) In particular, the Python interpreter on the laptop in front of you is a Turing machine.

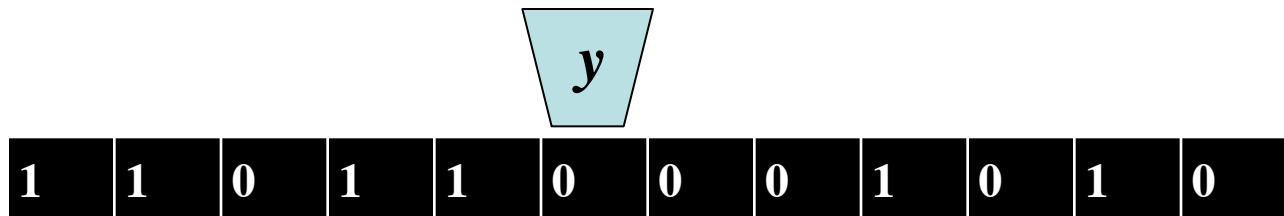


Turing Machines

- 1) The standard model of computation (up to and including human “computers”)
- 2) Almost all binary-valued functions $f(\cdot)$ over bit strings cannot be computed by any TM.

Proof: Set of all TMs = $\{0, 1\}^*$, the set of all finite bit strings

Set of all $f(\cdot) = 2^{\{0, 1\}^*}$. **QED.**



Smallest input to a Turing Machine

Kolmogorov complexity of bit string v , $K(v)$:

Minimum number of non-zero bits in an input bit string that causes the Turing machine to produce output string v and halt.

- Very common (and powerful) measure of “how complex” v is.
- Related to Shannon entropy (“complexity” of a singleton rather than of a distribution)
- “*Uncomputable*”, i.e., no computer program can calculate it. (There is no Turing machine that takes any v as input and eventually produces $K(v)$ as output and halts.)

- Generate input strings s to a TM by *coin-flipping* distribution:

$$P(s) = 2^{-|s|} / Z \quad (\text{Normalization constant } Z \leq 1)$$

- **Kolmogorov complexity** of bit string v , $K(v)$:

Minimum *length* input string s to a given TM
for it to compute v and halt.

- Bayes theorem:

$K(v)$ is length of most probable input s , given that output = v

Set thermo. rev. input distribution $q_0(s)$ to coin-flipping distribution

Thermodynamic complexity of bit string v :

Minimum *heat flow* for any input distribution $p_0(s)$ to a TM (that is reversible for $q_0(s)$) to compute v and then halt:

$$K(v) + \log[G(v)] + \log[Z]$$

where

- $K(v)$ is Kolmogorov complexity of v
- Z – the normalization constant – is Chaitin's constant
- $G(v)$ is probability of v under $q_0(s)$

Wolpert, D., *J. Phys. A* (2019)

Kolchinsky, A., Wolpert, D., *Phys. Rev. R* (2020)

Set thermo. rev. input distribution $q_0(s)$ to coin-flipping distribution

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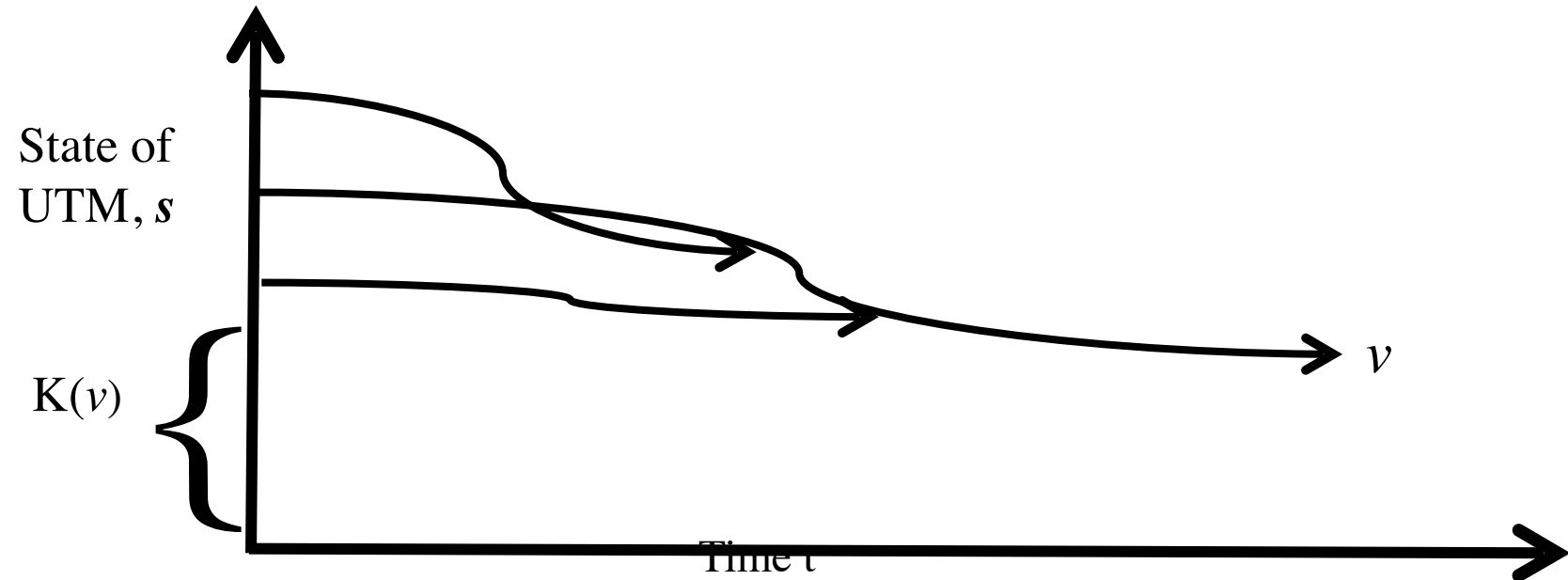
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where

- $K(v)$ is Kolmogorov complexity of v
- Z – the normalization constant – is Chaitin's constant
- $G(v)$ is probability of v under $q_0(s)$

A “correction” to Kolmogorov complexity, reflecting
cost of many-to-one maps as the TM evolves

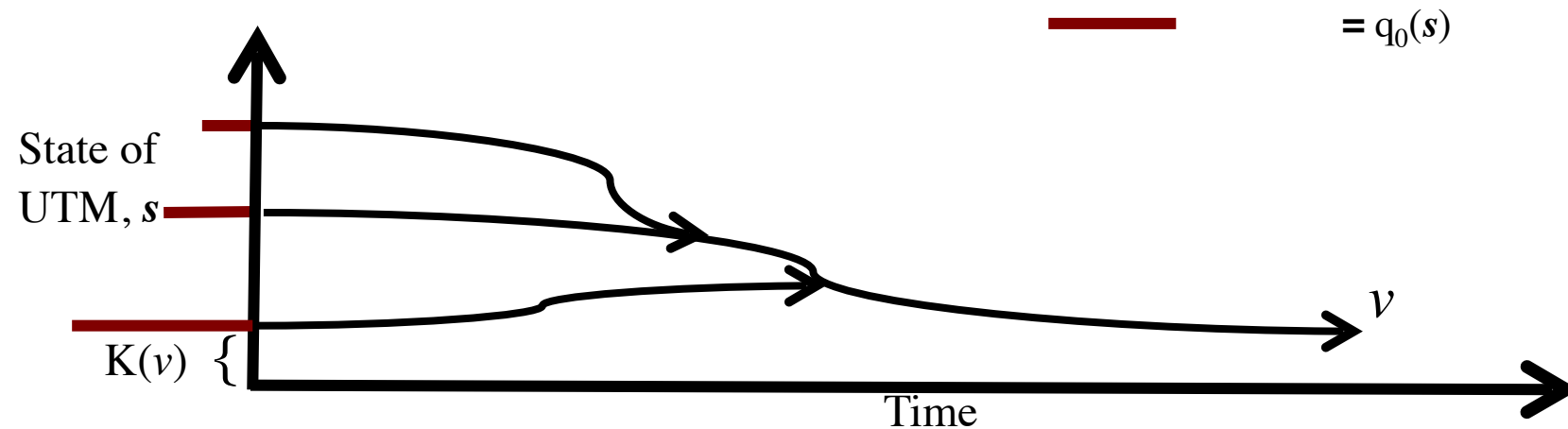
Kolmogorov complexity of v :



$K(v)$ is *unbounded* – no constant exceeds length of {the shortest string to compute v } for all v

Thermodynamic complexity of v :

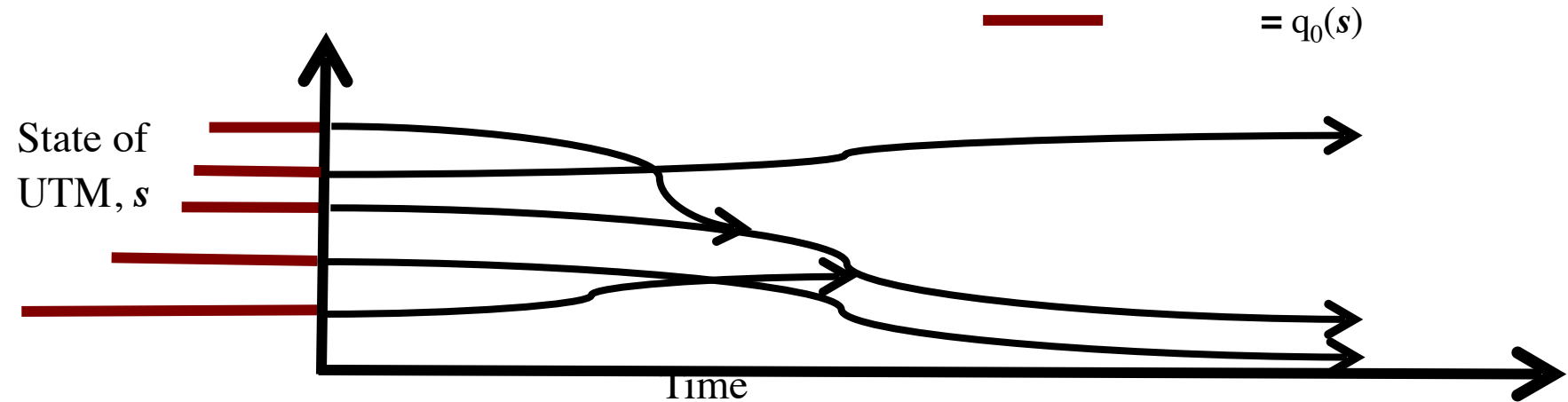
$$K(v) + \log[G(v)] + \log[Z]$$



Minimal heat flow is **bounded** – there is a constant that exceeds {minimal EF to compute v } for all v :
 $\log[\text{sum of lengths of red lines}]$

Thermodynamic complexity of v :

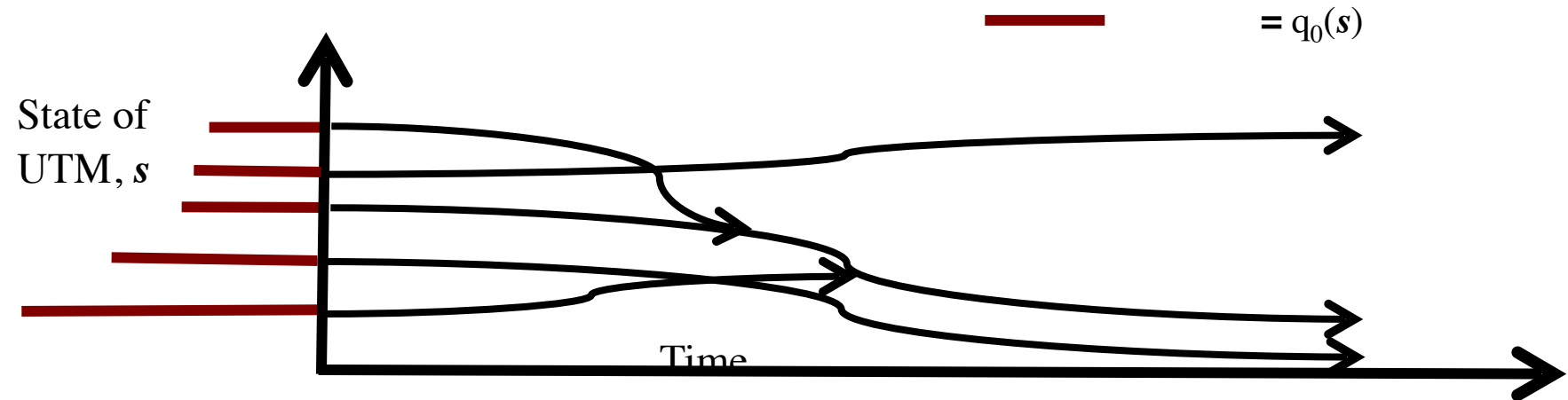
$$K(v) + \log[G(v)] + \log[Z]$$



Expected heat flow is infinite

Thermodynamic complexity of v :

$$K(v) + \log[G(v)] + \log[Z]$$



How do these results get modified if the update function of the TM is thermodynamically inefficient?

Answer: Who knows?

(See papers by Brittain et al., and Strasberg et al., in bibliography)

SUMMARY OF RESULTS

- Equation for heat expelled by a dynamic system (like a computer):

$$EF(p_0) = \text{Landauer cost}(p_0) + EP(p_0)$$

- Now have broadly applicable bounds on the second resource cost, $EP(p_0)$**
- Here consider two such bounds on that resource cost, **SLT** and **mismatch cost**.

-
- Both SLT and mismatch cost distinguish among computationally equivalent circuits with identical number of gates
 - Rich behavior of mismatch cost for DFAs run for multiple iterations
 - Nontrivial scaling of mismatch cost with input size for insert-sort algorithm
 - Thermodynamic Kolmogorov complexity is **bounded** (unlike Kolmogorov complexity)
 - Average thermodynamic work to run a TM is **infinite**

BIBLIOGRAPHY (a bit out of date)

Analysis of **electronic components** used in digital computers

- Riechers, P., in “*The Energetics of Computing in Life and Machines*”, Wolpert, D. et al. (Ed.’s), SFI Press (2019)
- Freitas, N., Delvenne, J., Esposito, M., *arxiv:2008.10578* (2021)
- Gao, C., Limmer, D., *arxiv:2102.13067* (2021)
- Boyd, A. Riechers, P., Wimsatt, G., Crutchfield, J., Gu, M., *arxiv:2104.12072* (2021)

Analysis of **Turing machines** - *based on stochastic thermodynamics*

- Strasberg, P., Cerrillo, J., Schaller, G., Brandes, T., *Phys. Rev. E* (2015)
- Wolpert, D., *J. Phys. A* (2019)
- Kolchinsky, A., Wolpert, D., *Phys. Rev. R* (2020)
- Brittain, R., Jones, N., Ouldrige, T., *arxiv:2102.03388*

Analysis of **Turing machines** - *not based on stochastic thermodynamics*

- Zurek, W., *Phys. Rev. A* (1989)
- translation into stochastic thermodynamics in Wolpert, D., *J. Phys. A* (2019)
- Bennett, C., *IBM J. Res. Dev.* (1973)

BIBLIOGRAPHY (a bit out of date)

Analysis of **straight-line programs** (including Bayes nets)

- Ito, S., Sagawa, T., *Phys. Rev. Letters* (2013)
- Ito, S., Sagawa, T., in “*Mathematical Foundations and Applications of Graph Theory*”, Dehmer M., et al. (Ed.’s), Wiley (2015)
- Wolpert, D., *J. Phys. A* (2019)
- Wolpert, D., Kolchinsky, A., *New J. Phys.* (2020)
- Wolpert, D., *Phys. Rev. Letters* (2020)

Analysis of **finite state automata** (including Mealy machines)

- Ganesh N., Anderson N., *Phys. Lett. A* (2013)
- Chu D., Spinney R., arXiv:1806.04875 (2018)
- Garner A., Thompson J., Vedral V., Gu M., *Phys. Rev. E* (2017)
- Boyd A., Mandal D., Crutchfield J., *New J. Phys.* (2016)
- Boyd A., Mandal D., Crutchfield J., *Phys. Rev. E* (2017)
- Boyd A., Mandal D., Riechers P., Crutchfield J., *Phys. Rev. Lett.* (2017)
- Boyd A., Mandal D., Crutchfield J., *J. Stat. Phys.* (2017)

BIBLIOGRAPHY (a bit out of date)

Analysis of arbitrary asynchronous information processing systems

- Sagawa, T., Ueda, M., *Phys. Rev. Letters* (2009)
- Sagawa, T., Ueda, M., *Phys. Rev. Letters* (2012)
- Sagawa, T., Ueda, M., *New. J. Phys.* (2013)
- Horowitz, J., Esposito, M., *Phys. Rev. X* (2014)
- Barato, A., Hartich, D., Seifert, U., *New. J. Phys.* (2014)
- Horowitz, J., *J. Stat. Mech.: Th. and Exp.* (2015)
- Barato, A., Seifert, U., *New. J. Phys.* (2017)
- Hartich, D., Barato, A., Seifert, U., *Phys. Rev. E* (2016)
- Brittain, R., Jones, N., Ouldrige, T., *J. Stat. Mech.: Th. and Exp.* (2017)
- Kardeş, G., Wolpert, D., *arxiv:2102:01610* (2020)
- Wolpert, D., *arxiv:2003:11144* (2020)
- Wolpert, D., *New J. Phys.* (2020)

BIBLIOGRAPHY (a bit out of date)

Thermodynamic (Ir)relevance of logical reversibility

- Maroney, O., *Phys. Rev. E* (2009)
- Sagawa, T., *J. Stat. Mech.: Th. and Exp.* (2014)
- Wolpert, D., *J. Phys. A* (2019)

Miscellaneous

- Parrondo., J., Horowitz, J., Sagawa, T., *Nature Physics* (2015)
- Sheng, S., Herpich, T. Diana, G., Esposito, M., *Entropy* (2019)
- Wolpert, D., Kempes, C., Stadler, P., Grochow, J., “*The Energetics of Computation in Life and Machines*”, SFI Press (2018)
- Grochow, J., Wolpert, D., *ACM SIGACT News* (2018)
- Wolpert, D., Kolchinsky, A., Owen, J., *Nat. Comm.* (2019)
- Owen, J., Kolchinsky, A., Wolpert, D., *New J. Phys.* (2019)
- Riechers, P., Gu, M., *Phys. Rev. E* (2021)
- Kolchinsky, A., Wolpert D., *arxiv:2103.05734*

BIBLIOGRAPHY (a bit out of date)

Relevant **computer science**; **comp. in biology**; **comp. in foundations of physics** (very partial)

- Nielsen, M., *Phys. Rev. Letters* (1997)
- Pour-El, M., Richards, J., “*Computability in Analysis and Physics*” (1997)
- Tegmark, M., *Annals of Phys.* (1998)
- Hut, P., Alford, M., Tegmark, M., *Found. Phys.* (2006)
- Li, M., Vitanyi, P., “An introduction to Kolmogorov Complexity and its applications”, Springer (2008)
- Soloveichik D., Cook M., Winfree E., and Bruck J., *Nat. Computing* (2008)
- Tegmark, M., *Found. Phys.* (2008)
- Arora, S., Barak, B., “*Computational Complexity: A modern approach*”, CUP (2009)
- Prohoska, S., Stadler, P., Krakauer, D., *J. Theor. Bio.* (2010)
- Barrow, J., in “*Kurt Godel and the Foundations of Mathematics*”, Baaz M., et al. (Ed.’s), CUP (2011)
- Aaronson, S., in “*Electronic Colloquium on Computational Complexity*” (2011)
- Qian, L, Winfree, E., *Science* (2011)
- Benenson, Y., *Nat. Rev. Genetics* (2012)
- Cubitt, T., Garcia-Perez, D., Wolf, M., *Nature* (2015)
- Allaghi, A., Hayes, J., *IEEE Trans. CAD of ICs and Systems* (2015)
- Barua, B., Mondal, C.; *J. Inst. Engineers: Series B*, (2019)
- Shiraishi, N., Matsumoto, K., *arxiv:2012.13890*